# Probing the nature of DM with CMB observations

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#### The CMB brightness is 2-dimensional function: frequency and direction

 $T(\nu, \hat{n}) = T_0 + \Delta T_{\text{spec}}(\nu) + \Delta T_{\text{anis}}(\hat{n}) + \delta T(\nu, \hat{n})$ 





**Planck** polarization



Anisotropies tell us about initial conditions + dynamics of perturbations + ionization history

$$\frac{\Delta T}{T}(z=0,\hat{n}) = \int dz \mathcal{V}(z) \mathcal{T}(z,r(z)\hat{n}) \zeta_{\text{init}} [+ (I)SW]$$

**Transfer function** (mostly photon monopole + dipole for temperature, quadrupole for polarization): result of linear evolution of photon + neutrino + baryon + DM.

Visibility function: probability of last Thomson scattering between z and z + dz.

$$\mathcal{V}(z) = n_{\mathrm{H}} x_{e} \sigma_{\mathrm{T}} \frac{dt}{dz} \times \exp\left[-\int_{0}^{z} dz' \left(n_{\mathrm{H}} x_{e} \sigma_{\mathrm{T}} \frac{dt}{dz}\right)_{z'}\right]$$



0-th order measurement: total abundance of cold, collisionless dark matter

$$\Omega_c h^2 = 0.1199 \pm 0.0022$$

[Planck collaboration 2015]



What can the CMB tell us about the nature of DM?

## Testing the vanilla WIMP with the CMB

"WIMP miracle": DM abundance results from annihilation cross section

$$\langle \sigma v \rangle_{\rm relic} \sim 10^{-26} {\rm cm}^3 {\rm s}^{-1}$$

#### **Step 0: DM** annihilations inject energy with volumetric rate

$$\dot{\rho}_{\rm inj} = \frac{1}{2} \left(\frac{\rho_c}{m_{\chi}}\right)^2 \langle \sigma v \rangle m_{\chi} = \frac{1}{2} \frac{\langle \sigma v \rangle}{m_{\chi}} \rho_c^2$$

 $p_{\rm ann} = \frac{\langle \sigma v \rangle}{m_{\gamma}}$ 

=> CMB is mostly sensitive to

#### Step 1: compute energy deposited into the plasma

$$\dot{\rho}_{\rm dep}(z) = f(z) \ \dot{\rho}_{\rm inj}$$

Depends on the nature and spectrum of particles produced in annihilation, e.g. if the DM annihilates entirely to neutrinos, f = 0



f(z) from Slatyer, Padmanabhan & Finkbeiner 09 for different models

#### Step 2: channel of energy deposition: heat / ionizations / excitations

Simple estimate of Chen & Kamionkowski 04 (based on numerical studies of Shull & Van Steenberg 85):

fraction in heat, ionization, excitation ~

$$\left[\frac{1+2x_e}{3}, \frac{1-x_e}{3}, \frac{1-x_e}{3}\right]$$

In practice, depends on detailed injection spectrum



#### **Step 3: effect on CMB observables**

• spectral distortions: 
$$\frac{\Delta I_{\nu}}{I_{\nu}^{\text{BB}}} \sim \int dt \ f_{\text{heat}} \frac{\dot{\rho}_{\text{dep}}}{\rho_{\gamma}}$$

typically much less sensitive to DM annihilation than anisotorpies

•recombination history: 
$$\dot{x}_e^{\text{direct}} = \dot{x}_e^{\text{std}} + f_{\text{ion}} \frac{\dot{\rho}_{\text{dep}}}{n_{\text{H}} \times 13.6 \text{ eV}}$$

$$\dot{x}_2 = \dot{x}_2^{\text{std}} + f_{\text{exc}} \frac{\dot{\rho}_{\text{dep}}}{n_{\text{H}} \times 10.2 \text{ eV}}$$

$$\dot{T}_{\rm gas} = \dot{T}_{\rm gas}^{\rm std} + \frac{2}{3n_{\rm gas}} f_{\rm heat} \ \dot{\rho}_{\rm dep}$$

Requires highly accurate standard recombination theory [YAH & Hirata 10, 11, Chluba ++ 11]

#### Implemented by several authors, e.g. Galli ++ 09, Slatyer ++09, Finkbeiner ++12, Green, Meerburg & Meyers 18



Giesen, Lesgourgues, Audren & YAH 12, computed with HyRec [YAH & Hirata, 2010, 2011]



Giesen, Lesgourgues, Audren & YAH 12, computed with HyRec + CLASS [Lesgourgues 2011]



#### Anisotropies are mostly sensitive to DM annihilation around z~400-800

weighing function f(z) derived by PCA [Slatyer 2016]

Same general picture holds to test any DM model injecting energy into the photon-baryon plasma (e.g. decaying DM):

energy injection -> energy deposition -> heat + ionizations

### CMB constraints on primordial black holes



### CMB constraints on primordial black holes

Ricotti, Ostriker & Mack 08, YAH & Kamionkowski 2017

Same idea: energy injection -> energy deposition -> heat + ionizations

accretion rate and radiative efficiency

mostly through Compton scattering of ~MeV photons







## Bounds to DM through effect on recombination history



## CMB as a direct detection experiment: testing dark matter scattering

suppose the DM elastically scatters with baryons / electrons / photons.

=> exchanges momentum with the photon-baryon plasma

=> affects the linear evolution of perturbations [Chen ++ 02, Sigurdson ++ 04, Dvorkin ++ 14, Boddy & Gluscevic 18]

$$a^{-1}\frac{d}{dt}(a\vec{v}_b) = -\vec{\nabla}\phi + \Gamma_{\text{Compt}}(\vec{v}_{\gamma} - \vec{v}_b) + \Gamma_{\chi b}(\vec{v}_{\chi} - \vec{v}_b)$$
$$a^{-1}\frac{d}{dt}(a\vec{v}_{\chi}) = -\vec{\nabla}\phi + \frac{\rho_b}{\rho_{\chi}}\Gamma_{\chi b}(\vec{v}_b - \vec{v}_{\chi})$$

$$\Gamma_{\chi b} \sim \frac{\rho_{\chi}}{m_b + m_{\chi}} \langle \sigma v_{\rm rel} \rangle$$

## Effect on CMB anisotropies



See also Boddy & Gluscevic 18 for constraints to general operators.

Bounds appear weaker than direct detection limits but

- (i) can probe masses < GeV
- (ii) shielding puts a ceiling to detectable cross sections

## **Constraints from spectral distortions**

If non-relativistic, scattering DM extracts heat from the photon-baryon plasma

In progress: detailed study of thermal decoupling of DM and resulting spectral distortion bounds.



## Summary

- CMB anisotropies and spectral distortions can constrain DM annihilation, or in general energy injection in the plasma, e.g. PBHs.
- The CMB can also probe scattering DM through its cooling effect and momentum exchange.
- How much can bounds be improved?



Figure 4. Improvement on forecasted constraints for a cosmic-variance-limited temperature-only CMB experiment up to a given  $\ell_{\text{max}}$  relative to that with  $\ell_{\text{max}} = 1500$ .