Radiogenic Neutron Background v1.0

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Ultra-low-background experiments, generically termed rare-event searches, address some of the most important open questions in particle physics, cosmology and astrophysics: direct detection of dark matter, neutrinoless double beta decay, proton decay, and detection of solar and supernovae neutrinos. Although their detection methods and physics goals are varied, rare-event searches share a number of common requirements in order to obtain sensitivity to low-rate processes in the presence of an overwhelming rate of environmental radiation, including the need for significant rock overburdens to moderate the flux of cosmic-ray muons. This requirement is so generic that an international community of underground science has emerged, with dozens of experiments operating in mines that vary in depth from a few hundred meters water equivalent (m.w.e.) to many thousand.

Simulations are used to understand backgrounds caused by naturally occurring radioactivity in the rock and in every piece of shielding and detector material used in the experiment. Most important are processes like spontaneous fission and (α,n) reactions in material close to the detectors that can produce neutrons. A comparison study between two dedicated softwares is detailed. Neutron yields and spectra obtained with Mei-Zhang-Hime and SOURCES4 codes are presented.

I. INTRODUCTION

Reduction of radiogenic backgrounds is one of the most important factors for rare event search experiments including searches for dark matter and neutrino-less double beta decay. These radiogenic backgrounds can be generically classified into two types: electron recoil and nuclear recoil backgrounds. Electron recoil background result from the interactions of gamma rays, electrons or beta particles interacting with the electrons in the detector's target medium while nuclear recoil backgrounds result from neutrons interacting with the nucleus in the detector's target medium. In dark matter experiments, the nuclear recoil background is of particular concern because a nuclear recoil from a neutron can be indistinguishable from a nuclear recoil from a WIMP.

Neutrons are produced by spontaneous fission, α-n interactions and muon-induced interactions. Muon-induced neutrons can be reduced by operation the detector deep underground and by placing passive shielding around the detector. In addition, muon-induced neutrons can often be recognized by identifying the parent muon in a muon veto. More difficult to deal with are neutrons resulting from α-n interactions and spontaneous fission from ^{238}U, ^{235}U, and ^{232}Th present in the materials the shielding and the detectors themselves. Thus, experimentalists must select their materials carefully when designing and constructing their experiments. This requires that they have accurate simulations of the processes in the materials they are considering.

There are currently two codes available for use in such calculations, SOURCES-4 and a second code developed by Mei, Zhang and Hei which we will refer to as the USD code.

II. NEUTRON YIELD AND SPECTRA

The measurements of neutron spectra strongly depend on the material and are not straightforward since neutrons are neutral particles. Their calculations is critical to low background experiments such as direct dark matter search experiments. The total neutron yield indicates the number of neutrons which are produced or had entered the target whereas the neutron energy spectrum determines the background events we would expect in the energy range of interest. Both are then needed in order to carry out a complete and reliable neutron background simulation. The evaluation of neutron yields and spectra can be performed via different codes: the modified version of SOURCES4A [1] and the code developed by Mei, Zhang and Hime [2] made available online at http://neutronyield.usd.edu have been considered. SOURCES4 code has been modified to extend the cross section for (α,n) up to 10MeV, based on experimental data, whenever possible, and calculations performed via the EMPIRE code [3] by the group at the University of Sheffield. SOURCES4A and USD codes calculate neutron yields and spectra from (α, n) reactions due to the decay of radionuclides. Radiogenic neutrons are due to intrinsic contamination of materials surrounding the detectors of ^{232}Th, ^{238}U and ^{235}U decay chains. Comparison study between SOURCES4 and USD codes is carried out considering ^{232}Th, ^{238}U decay chains in secular equilibrium, although a possibility of disequilibrium can be taken into account: due to migration differently, the long-lived isotopes, ^{226}Ra, ^{222}Rn, ^{210}Po, ^{228}Ra, ^{228}Th and the associated decay daughters could be calculated separately, see table I. For both chains, most of the neutrons are produced by the α generated in the second part of the chains.

TABLE I. Radiogenic neutron yield (n/s/cm^3) from (α, n) reactions in different materials for 1ppb of ^{238}U and ^{232}Th decay chains. Neutron yield have been calculated via the modified SOURCES4A code.

	Neutron yield for 1ppb (n/s/cm^3)	
Material	**^{232}Th $\rightarrow$ ^{228}Th**	**^{228}Th $\rightarrow$ ^{208}Pb**
Stainless Steel	$8.3 \cdot 10^{-15}$	$4.1 \cdot 10^{-11}$
Pyrex	$1.1 \cdot 10^{-11}$	$8.4 \cdot 10^{-11}$
Hamamatsu PMT	$1.7 \cdot 10^{-11}$	$1.2 \cdot 10^{-10}$
Titanium	$2.0 \cdot 10^{-13}$	$9.3 \cdot 10^{-11}$
Copper	$0.0 \cdot 10^{-11}$	$9.5 \cdot 10^{-12}$
PE (CH_2)	$4.9 \cdot 10^{-13}$	$5.1 \cdot 10^{-12}$
PTFE (CF_2)	$6.3 \cdot 10^{-11}$	$7.7 \cdot 10^{-10}$
	^{238}U $\rightarrow$ ^{226}Ra	**^{226}Ra $\rightarrow$ ^{206}Pb**
Stainless Steel	$1.0 \cdot 10^{-14}$	$3.1 \cdot 10^{-11}$
Pyrex	$5.6 \cdot 10^{-11}$	$1.9 \cdot 10^{-10}$
Hamamatsu PMT	$8.9 \cdot 10^{-11}$	$2.8 \cdot 10^{-10}$
Titanium	$1.8 \cdot 10^{-13}$	$1.0 \cdot 10^{-10}$
Copper	$0.0 \cdot 10^{-11}$	$2.8 \cdot 10^{-12}$
PE (CH_2)	$2.2 \cdot 10^{-12}$	$1.1 \cdot 10^{-11}$
PTFE (CF_2)	$2.6 \cdot 10^{-10}$	$1.6 \cdot 10^{-9}$

The calculation of the neutron spectra requires as input the cross-sections of (α, n) reactions, the probabilities of nuclear transition to different excited states (branching ratio) and the alpha emission lines from the radioactive radionuclides. Both codes consider a thick target: calculation of neutron yields and spectra are carried out under the assumption that the size of radioactive sample exceeds significantly the range of alpha. The energy bin size of the (α, n) calculation is fixed to be 0.1 MeV in USD code while is user dependent in SOURCES4A.

Table II lists α-lines respectively for ^{232}Th and ^{238}U decay chains, present in SOURCES4A and USD codes. USD code is missing the lines coming from ^{222}Ra isotope in ^{238}U decay chain which SOURCES4A library is considering. Overall, the branching ratio and the energy lines are in good agreement between the two codes.

Cross-sections and branching ratios are required as well for neutron yield and spectra calculations. The USD code considers TENDL 2012 [4] to provide (α, n) nuclear cross-sections. TENDL is a validated nuclear data library which provides the output of the TALYS [5] nuclear model code system; SOURCES4 cross section input libraries come from EMPIRE2.19 [3] calculations and, for some isotopes, a combination of data measurements and EMPIRE2.19 calculations. Also, EMPIRE is the code recommended by International Atomic Energy Agency (IAEA). Neither EMPIRE nor TALYS can calculate properly resonance behavior which has been experimentally observed.

TABLE II. Alpha lines present in SOURCES4A and USD, and their intensity (BR) for isotopes in the ^{232}Th and ^{238}U decay chains. Only lines with intensity > 1% have been quoted.

	SOURCES4A		**USD**	
Isotopes	**Line (keV)**	**BR (%)**	**Line (MeV)**	**BR (%)**
^{238}U	4151	21	4.151	21
	4198	79	4.198	79
^{234}U	4722.4	28.42	4.722	28.6
	4774.6	71.38	4.775	71.4
^{230}Th	4620.5	23.4	4.621	23.7
	4687.0	76.3	4.688	76.3
^{226}Ra	4601	5.55	4.602	5.6
	4784.34	94.45	4.784	94.4
^{222}Ra	6241	3.05	na	na
	6559	96.90	na	na
^{218}Po	6002.35	99.98	6.002	100
^{214}Po	7686.82	99.99	7.687	99.99
^{210}Po	5304.33	99.99	5.304	100
^{232}Th	3947.2	21.7	3.954	22.1
	4012.3	78.2	4.013	77.9
^{228}Th	5340.36	27.2	5.340	28.5
	5423.15	72.2	5.423	71.5
^{224}Ra	5448.6	5.06	5.449	5.1
	5685.37	94.92	5.685	94.9
^{220}Ra	6288.3	99.99	6.288	100
^{216}Po	6778.5	99.99	6.778	100
^{212}Bi	6051.1	25.16	6.050	26.2
	6090.2	9.79	6.0902	9.8
^{212}Po	8784.6	100	8.784	64

TABLE III. Radiogenic neutron yield (n/s/cm^3) for copper and polyethylene materials and for ^{238}U and ^{232}Th decay chains. Column (1) and (2) refer to pure USD and SOURCES4A calculation, respectively. Column (3) refers to SOURCES4A calculation with USD (α, n) cross section libraries. A ratio of the neutron yield is also provided: column (a) refers to the ratio of (2) over (1), whereas column (b) corresponds to the ratio (2)/(3)

		Neutron Yield (n/s/cm^3)			**Ratio**	
Material	**Chain**	**(1)**	**(2)**	**(3)**	**(a)**	**(b)**
Cu	^{238}U	3.46E-12	2.84E-12	2.93E-12	0.8	1.0
	^{232}Th	1.11E-11	9.49E-12	9.18E-12	0.9	1.0
PE	^{238}U	9.56E-12	1.26E-11	1.64E-11	1.3	0.8
	^{232}Th	2.87E-12	5.28E-12	5.97E-12	1.8	0.9

An extensive comparison between the different cross-section libraries used in USD and SOURCES4A has been carried out for the target nuclides present in the materials contributing the most to radiogenic neutron background. Specifically, we refer to materials which compose the shielding scheme and the internal detector parts, such as stainless steel, copper, titanium, borosilicate glass (PMTs), PTFE which become important when the external flux is attenuated by the shielding. Comparison of cross section inputs results in a good agreement for most of isotopes in both code libraries. For some, such as ^{13}C and ^{10}B and ^{11}B the cross sections show discrep-

FIG. 1. Input (α, n) cross-section for the target isotopes involved in radiogenic neutron calculations for copper and polyethylene. Red markers refers to SOURCES4A inputs, whereas the black markers to USD. From left to right: ^{13}C, ^{63}Cu and ^{65}Cu

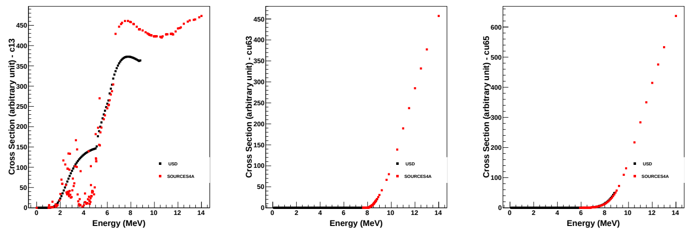
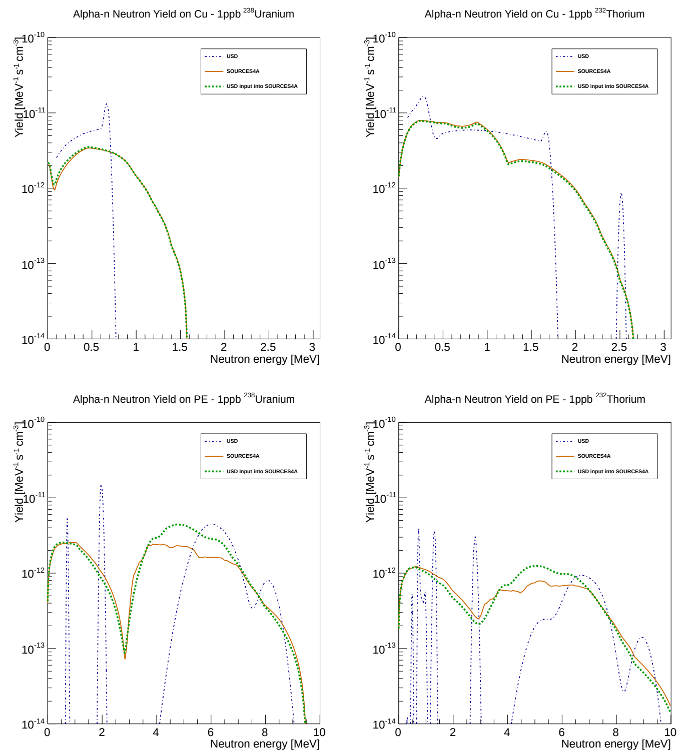


FIG. 2. Radiogenic neutron spectra (n/MeV/s/cm^3) calculated for 1ppb ^{238}U and ^{232}Th decay chains, left and right panels, respectively. First row show copper contribution, lower row polyethylene material. Dotted blu lines refers to pure USD calculations, plain orange line to pure SOURCES4A calculations and dashed green line is the mixed computation for which we ran SOURCES4A algorithm with input cross section of USD code.



ancies. To better understand the contribution due to the cross section input library in the radiogenic neutron yield and spectrum we have calculated radiogenic neutron yield and spectra for two different materials: copper, for which the input cross sections in both codes are matching, and polyethylene, for which input cross sections of ^{13}C show discrepancies between the codes, details in figure 1. We have considered natural copper (70% ^{63}Cu and 30% ^{65}Cu) with a density of 8.96 g/cm^3; polyethylene material (CH_2) is considered with a density of density is 0.935 g/cm^3. Estimates are done for 1ppb in ^{238}U and ^{232}Th decay chains. Calculations consist of pure SOURCES4A and USD computation and a mixed computation for which we run SOURCES4A algorithm with cross section inputs of USD. Resulting radiogenic neutron yields are quoted in table III. The column (a) of the radiogenic neutron yield ratio quotes the calculation differences between SOURCES4A and USD codes: they show a reasonably good agreement, inside a 50% discrepancy. The column (b) refers to the ratio of radiogenic neutron yield resulting from the same algorithm calculation (SOURCES4A) but with USD (column (3)) and SOURCES4A (column (2)) input (α, n) cross-sections. For polyethylene, we can conclude that the input cross-section may account up to a 20% discrepancy in neutron yield. Figure 4 shows radiogenic neutron spectra for copper (upper row) and polyethylene (lower row), both from uranium and thorium decay chains, left and right panels respectively.

A comparison of neutron yield and energy spectrum obtained via SOURCES4A and USD codes for different material has been carried out. Results are shown in table IV and figure 4. A qualitative agreement between the two codes is observed. A maximum discrepancy of a factor 2 is found. The energy spectra calculated via SOURCES4A code are in general smoother, without the presence of peaks, feature of USD spectra.

TABLE IV. Radiogenic neutron yield (n/s/cm^3) per material considering 1ppb of ^{238}U and ^{232}Th.

		Neutron Yield (n/s/cm^3)			
Material	**Chain**	SOURCES4A	USD	Ratio	Diff %
Copper	^{238}U	3.46E-12	2.84E-12	1.2	20
	^{232}Th	1.11E-11	9.49E-12	1.2	16
PE (CH_2)	^{238}U	9.56E-12	1.26E-11	0.8	-27
	^{232}Th	2.87E-12	5.28E-12	0.5	-59
Titanium	^{238}U	1.04E-10	1.99E-10	0.5	-63
	^{232}Th	9.29E-11	1.24E-10	0.8	-28
Stainless Steel	^{238}U	3.10E-11	5.95E-11	0.5	-63
	^{232}Th	4.05E-11	6.80E-11	0.6	-51
Pyrex	^{238}U	2.30E-10	1.61E-10	1.4	36
	^{232}Th	8.66E-11	4.59E-11	1.9	61
Hamamtsu PMT	^{238}U	3.48E-10	2.45E-10	1.4	35
	^{232}Th	1.27E-10	6.98E-11	1.8	58
PTFE (CF_2)	^{238}U	1.81E-09	1.60E-09	1.1	12
	^{232}Th	7.76E-10	5.42E-10	1.4	36

III. GEANT4 PROPAGATION

Simple MC propagation of spectra to check how the discrepancies from USD website and SOURCES4 calculations can affect experimental simulations. It has been done for Cu using Geant4.9.5 considering a single Cu can (21760.6 cm3) around 100kg of germanium detector. Borosilicate glass neutron spectrum has been propagated within RAT, utilizing Geant4.9.5, and a cylindrical 45T liquid argon single phase detector surrounded by borosilicate glass mimicking PMTs as well as stainless steel and a water veto.

other simulations?

IV. CONCLUSION

[1] W. Wilson, R. Perry, W. Charlton, and T. Parish, Progress in Nuclear Energy **51**, 608 (2009).

[2] D. M. Mei, C. Zhang, and A. Hime, Nucl. Instrum. Meth. **A606**, 651 (2009), arXiv:0812.4307 [nucl-ex].

[3] E.-I. statistical model code for nucelar reaction calculation, https://www-nds.iaea.org/empire218/manual.ps..

[4] A. Koning, D. Rochman, S. van der Marck, J. Kopecky, J. C. Sublet, S. Pomp, H. Sjostrand, R. Forrest, E. Bauge, and H. Henriksson, TENDL-2012: TALYS-based evaluated nuclear data library, ftp://ftp.nrg.eu/pub/www/talys/tendl2012/tendl2012.html.

[5] A. Koning and D. Rochman, Nuclear Research and Consultancy Group (NRG), http://www.talys.eu/documentation/.

FIG. 3. Radiogenic neutron spectra (n/MeV/s/cm^3) calculated for 1ppb ^{238}U and ^{232}Th decay chains, left and right panels, respectively. The (α, n) reaction contribution is shown in variuos commonly used materials from SOURCES4A in orange and USD in blue. From top to bottom materials are: copper, titanium, stainless steel, pyrex, borosilicate glass (Hamamtsu PMT galss), polyethylene and teflon (PTFE).

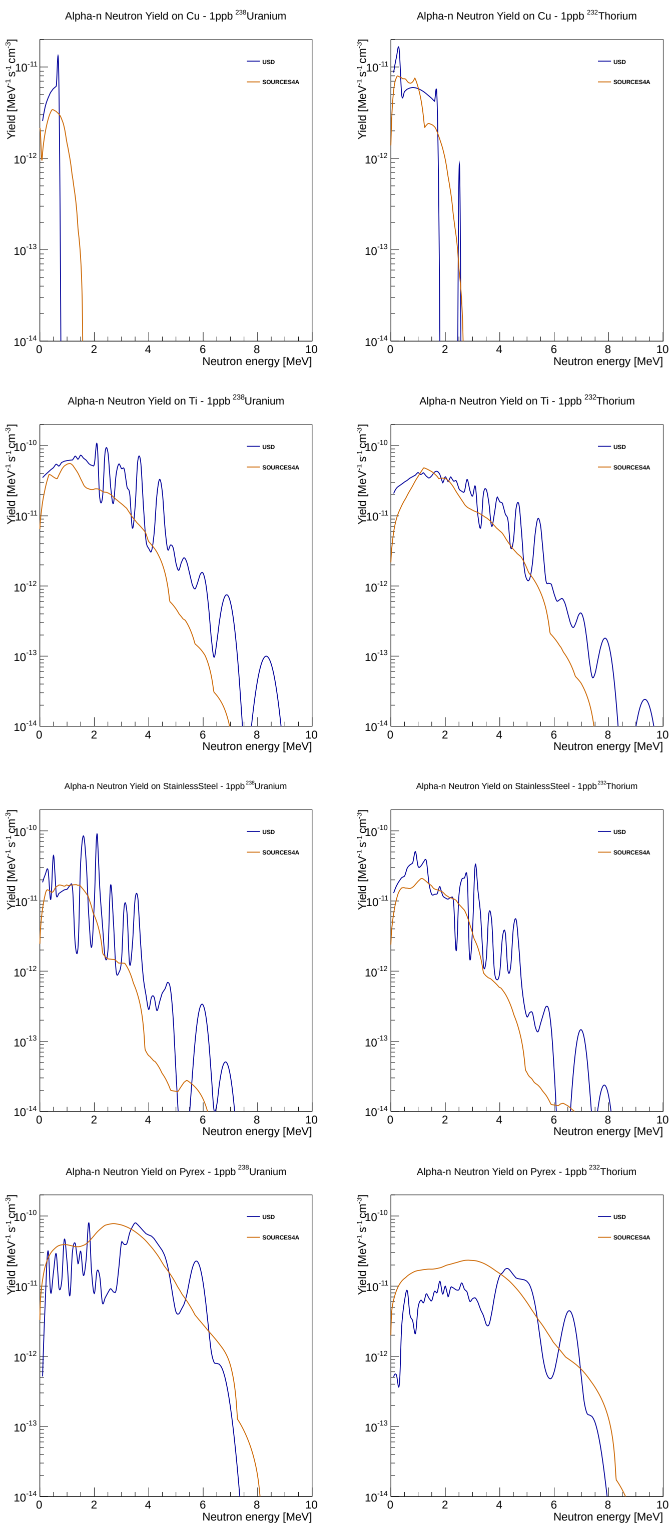


FIG. 4. Radiogenic neutron spectra (n/MeV/s/cm^3) calculated for 1ppb ^{238}U and ^{232}Th decay chains, left and right panels, respectively. The (α, n) reaction contribution is shown in variuos commonly used materials from SOURCES4A in orange and USD in blue. From top to bottom materials are: copper, titanium, stainless steel, pyrex, borosilicate glass (Hamamtsu PMT galss), polyethylene and teflon (PTFE).