A Comprehensive Comparison for Simulations of Cosmic Ray Muons Underground

P. Cushman*, A. Empl[†], A. Kennedy^{*}, S. Lindsay[†] and A.N. Villano^{*}

*University of Minnesota, Minneapolis MN 55455 †University of Arkansas at Little Rock, Little Rock AR 72204

Abstract. The two leading simulation frameworks used for the simulation of cosmic ray muons underground are FLUKA and Geant4. There have been in the past various questions raised as to the equivalence of these codes regarding cosmogenically produced neutrons and radioactivity in an underground environment. Many experiments choose one of these frameworks and because they typically have different geometries and are located at different underground sites the issues relating to code comparison are compounded. We report on an effort to compare the results of each of these codes in simulations which have simple geometry which is consistent between the two codes. It is seen that in terms of a basic flux variable and neturon capture statistics the codes agree well in a broad sense. There are, however, differences that will be subject of further study. Comparisons of the simulations to available data are considered and the difficulties of such comparisons are pointed out.

Keywords: cosmogenics, neutrons, low background **PACS:** 96.50.S-,25.30.Mr,25.20.-x,25.30.-c

1. INTRODUCTION

Low background counting experiments have been very useful tools for providing information related to the physics of the standard model and beyond. Some of the most obvious examples are neutrino detection experiments, direct dark matter searches, searches for neutrinoless double-beta decay and searches for proton decay. Cosmogenically produced prompt radiation or residual radioactivity are sources of backgrounds for these experiments which can have impacts up to and beyond depths of 3 km.w.e if the experimental sensitivity is not limited first by some other background [1]. One approach which is typically taken to quantify these backgrounds is to propagate surface muon energy and angular distributions [2] to the depth of the experimental installation where a full microscopic simulation of the surviving muons in underground cavern material is performed []. Different experimental collaborations typically perform these simulations with various generalized simulation packages, FLUKA and Geant4 being two examples [3, 4, 5]. It is important to fully understand the physics included in these simulation packages, how their implementations differ, and how the key results compare to data. Because of these facts an effort to comprehensively characterize important physics for cosmogenically induced radiation with both FLUKA and Geant4 in a simple geometrical environment is underway and first results are reported here.

FLUKA is a particle physics Monte Carlo simulation package which traditionally has been applied to cosmogenic background problems. Its predictions are predominantly based on original and well-tested microscopic models. The physics models are fully

A Comprehensive Comparison for Simulations of Cosmic Ray Muons Underground June 17, 2013 1

integrated in the code and can not be modified by the user. Details about the implemented hadronic models relevant to the production of cosmogenic neutrons can be found in []. Low energy neutrons with kinetic energies < 20 MeV are treated in the multi-group approach in FLUKA which requires a careful interpretation of results at the single event level. FLUKA 2011.2.17 from December 2012 was used for the simulation with the FLUKA default setting PRECISIO(n). Photonuclear interactions were enabled through the user option PHOTONUC and a more detailed treatment of nuclear de-excitation was requested with the options EVAPORAT(ion) and COALESCE(nce). In addition, the treatment of nucleus-nucleus interaction was turned on for all energies via the option IONTRANS.

Geant4 is in general a microscopic particle tracking Monte Carlo code which has been predominantly used in high energy particle physics and radiation protection. The code attempts to explicitly include all relevant interactions to a simulation in a modular, userdefined way and then use the coded models and stored relative probabilities (usually cross sections) of each interaction in a given material to decide the course of the tracked particle at the next simulation step. The simulation of a particle proceeds on a microscopic basis (interaction-by-interaction) until an energy limit is reached when it can be absorbed into the material by depositing energy or decays. While the code includes well-tested microscopic models for many processes, details of the physics models used can in general be modified by the user and in principle the user has access to the entire source code of Geant4. The Shielding physics list is used for the Geant4 simulations in this study []. The Shielding list is a parameterization of many hadronic and leptonic models which previously had to be included one-by-one. One of the original uses for this list was underground or low background experiments and it includes high precision neutron transport physics. The version of Geant4 used in this study is currently Geant4.9.5, but the Shielding list has been available since Geant4.9.4.

2. GEOMETRY AND MATERIALS

In order to get the greatest access to the physics of the cosmogenic cascades a simple cylindrical geometry is used with five materials containing a wide range of nuclei and common detector materials. Each material is taken with a density-weighted thickness equal to $3200 \ g/cm^2$ and a radius of $10 \ m$ so that the captured neutrons can range out and the capture statistics can be allowed to probe the low energy diffusion properties of the material in each given simulation package. Liquid scintillator (C₉H₁2, density 0.887 g/cm³), water (density 0.997 g/cm³), calcium carbonate (CaCO₃, density 2.710 g/cm³), iron (density 7.874 g/cm³) and lead (density 11.342 g/cm³) are used. All isotopic abundances are those which appear naturally. In addition the primary muon or anti-muon energies of 30, 100, 280 GeV and 1 TeV are used in this study. In Sec. 3 we specialize to the 280 GeV energy setting for muons only and in Sec. 4 we specialize to the liquid scintillator material, similar to the one used in Borexino [] with muons only. These specializations are made to keep the length of this article tractable.

Figure 1 shows the geometrical setup for the liquid scintillator material C_9H_{12} . The fiducial region displayed in the center represents half of the full thickness in total and is the region over which capture statistics are gathered, whereas the thin planes

perpendicular to the axis are the integration planes for the integrated neutron flux statistics. The FLUKA simulations report the flux variables at the central plane while the Geant4 simulations report the average of the flux across four planes including the central plane and three evenly-spaced planes left of center. This scheme allows probing of the systematic difference in flux variables at planes after the full shower has developed.



FIGURE 1. Cylindrical geometry used for the simulations. A specific example using the liquid scintillator material C_9H_{12} is used. The density of this material is 0.887 g/cm^2 . The muon primaries are incident on the axis of the cylinder from large x-coordinate to small.

3. NEUTRON FLUX VS. ENERGY

The one-directional neutron flux integrated over the detector cross section perpendicular to the momentum of the muon primary is a good parameter to summarize the neutron production behavior over a wide range in neutron energies. This flux is plotted in Fig. 2, normalized to the number of generated primaries and the histogram bin widths. The flux compare favorably between the FLUKA and Geant4 simulations in general. The liquid scintillator and water seem to be within 30% over most of the energy range with the largest excursions near 100 MeV neutron energy. Other materials follow this general agreement in the region between 100 MeV and the highest energies plotted. Below 10 MeV CaCO₃ has several resonant structures which are well tracked and an indication of successful implementation in the nuclear physics regime (and similar cross sections). The case is similar for iron with the acception of a structure around 3-7 MeV which appears in the Geant4 simulation but not the FLUKA. The lead material has a similar structure in the same energy region not reproduced by FLUKA and furthermore Geant4 begins to register dramatically less neutrons than FLUKA 10 keV and 5 MeV. This is



FIGURE 2. Integrated neutron fluxes across perpendicular planes for all materials down to neutron energy of approximately 1 keV. Each material is scaled by a power of 10 for ploting purposes, lead by 10⁰, iron by 10⁻¹, CaCO₃ by 10⁻², water by 10⁻³ and liquid scintillator by 10⁻⁴. The FLUKA simulations are shown in blue shades while the Geant4 simulations are shown in yellow/orange shades. FLUKA simulations are integraded over a plane at the geometrical center of the geometry wheras Geant4 simulations are averaged over the central plane plus three evenly spaced planes left of the central plane after the shower has fully developed.

the largest discrepancy revealed so far and is not yet understood. Lead is, however, an important material for many low background experiments and so investigation of this discrepancy and correlation to data if possible is important.

4. THERMALIZATION AND MULTIPLICITY

In a simple geometry such as the one being utilized in this study, it is typically easy to keep track of the number of neutrons that are produced in a given event in total. Further, since the cylindrical material slabs have such a large lateral distance, it is likely that almost all of the produced neutrons will remain inside the detection volume. To this end Fig. 3 shows plots which analyze various properties of "captured" neturons inside the fiducialized detector volume. Here captured means the last tracked point of all neutrons with energies less than 3.1 eV. The energy 3.1 eV corresponds closely to a FLUKA low energy neutron group and is used for technical reasons. The capture time spectrum of Fig. 3a shows that FLUKA produces slightly less neutrons than Geant4 does and additionally the capture time constant for FLUKA is XXX whereas the time constant for Geant4 is YYY. The difference in these is interesting and indicates a slight difference in the low energy neutron transport physics, which warrants further study.



FIGURE 3. a) Capture time distributions for all neutrons which capture with energies less than 3.1 eV. b) Lateral distance distributions for all neutrons which capture with energies less than 3.1 eV. c) Multiplicity distributions counting all captures which take place with energies less than 3.1 eV.

The other distributions look qualitatively similar but there are several discrepancies. Firstly there is slight normalization discrepancy as it is often the case that the FLUKA curves are bounded below the Geant4 curves. There is also a discrepancy in the lateral distance distributions which appears to be worse toward lower distance and with lower energy muon primaries. Finally the multiplicity distributions show a large discrepancy often a factor of 2-4 in the multiplicity one bin. These discrepancies are can be linked and understanding their origins is important to vet the transport physics of each of the simulations.

Some experiments like Borexino can measure very analagous quantities at depth []. Some experimental conditions should be used to match simulation to experiment, like the amount of dead time after a muon traversal that the experiment will experience before being able to detect neutron captures. After these corrections, however, the capture time, lateral distance and multiplcity distributions should be quantitatively similar to the ones displayed in Fig. 3. Quantitative comparisons of that type have not been undertaken for these simulation data but will be in the future.

5. CONCLUSIONS

A typical way to summarize the quality of the agreement of data and simulations has been to plot the total neutron yields for a given material and energy []. The energy is usually quoted as the average energy of an underground muon spectrum but here in this simple study we use the energy of the muon primaries. Figure 4 displays this yield for the liquid scintillator material. Though the values track the Geant4 simulation produces 30-50% more neutrons in total than FLUKA. An understanding of these values could be related to the discrepancies noted in previous sections.

While there is good qualitative agreement in much of the simulation work presented



FIGURE 4. Normalized neutron yield for the liquid scintillator material at energies 30, 100, 280 GeV and 1 TeV primary muon energy. Geant4.9.5 produces about 30-50% more neturons on average with the settings used here.

here, there are discrepancies which need to be understood in terms of the computational physics being used and in terms of agreement with data. The analysis included here looks at some rough quantifiers of production (flux vs. energy) and transport (capture distributions) physics. The current analysis serves as the starting point for understanding the interplay of these types of variables and the investigation of new variables which can be utilized for benchmarking both between various Monte Carlo codes and between those codes and measured data. In the end the ideal reasult from these studies will be a detailed understanding of which observables constrain the important microscopic processes involved in cosmogenic studies, and a framework which can be used to benchmark any code which is put to use for deep underground cosmogenic simulations.

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A Comprehensive Comparison for Simulations of Cosmic Ray Muons Underground June 17, 2013 6

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